Aircraft millimeter-wave passive sensing of cloud liquid water and water vapor during VOCALS-REx

P. Zuidema1, D. Leon2, A. Pazmany3, and M. Cadeddu4

1Rosenstiel School of Marine and Atmospheric Sciences, University of Miami, Miami, Florida, USA
2Department of Atmospheric Sciences, University of Wyoming, Laramie, Wyoming, USA
3Prosensing Inc., Amherst, Massachusetts, USA
4Argonne National Laboratory, Argonne, Illinois, USA

Correspondence to: P. Zuidema (pzuidema@rsmas.miami.edu)

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Abstract. Routine liquid water path measurements and water vapor path are valuable for process studies of the cloudy marine boundary layer and for the assessment of large-scale models. The VOCALS Regional Experiment respected this goal by including a small, inexpensive, upward-pointing millimeter-wavelength passive radiometer on the fourteen research flights of the NCAR C-130 plane, the G-band (183 GHz) Vapor Radiometer (GVR). The radiometer permitted above-cloud retrievals of the free-tropospheric water vapor path (WVP). Retrieved free-tropospheric (above-cloud) water vapor paths possessed a strong longitudinal gradient, with off-shore values of one to two mm and near-coastal values reaching ten mm. The VOCALS-REx free troposphere was drier than that of previous years. Cloud liquid water paths (LWPs) were retrieved from the sub-cloud and cloudbase aircraft legs through a combination of the GVR, remotely-sensed cloud boundary information, and in-situ thermodynamic data. The absolute (between-leg) and relative (within-leg) accuracy of the LWP retrievals at 1 Hz (∼100 m) resolution was estimated at 20 g m$^{-2}$ and 3 g m$^{-2}$ respectively for well-mixed conditions, and 25 g m$^{-2}$ absolute uncertainty for decoupled conditions where the input WVP specification was more uncertain. Retrieved liquid water paths matched adiabatic values derived from coincident cloud thickness measurements exceedingly well. A significant contribution of the GVR dataset was the extended information on the thin clouds, with 62% (28%) of the retrieved LWPs <100 (40) g m$^{-2}$. Coastal LWPs values were lower than those offshore. For the four dedicated 20° S flights, the mean (median) coastal LWP was 67 (61) g m$^{-2}$, increasing to 166 (120) g m$^{-2}$ 1500 km offshore. The overall LWP cloud fraction from thirteen research flights was 63%, higher than that of adiabatic LWPs at 40%, but lower than the lidar-determined cloud cover of 85%, further testifying to the frequent occurrence of thin clouds.

1 Introduction

The VOCALS Regional Experiment, held in October–November 2008 in the southeast Pacific (SEP), had one focus on acquiring field measurements capable of illuminating the processes linking aerosols, clouds, and precipitation (Wood et al., 2011b), field measurements that could then help assess modeled cloud-aerosol interactions (e.g., Saide et al., 2011, and references therein). Cloud property datasets were also desired with which to assess large-scale climate models (Wood et al., 2011b). These considerations motivated the choice of instruments placed upon the NCAR C-130 plane. These included not only in-situ aerosol and microphysical probes, but also cloud remote sensors that could routinely provide information on stratocumulus precipitation and its vertical profile, the cloud boundaries, and liquid water (Wood et al., 2011b). These choices, given finite aircraft space, recognized the value of airborne remote sensors for cloudy boundary layer studies (e.g., Vali et al., 1998; van Zanten et al., 2005; Rauber et al., 2007; Bretherton et al., 2010; Wood et al., 2011a; Wang et al., 2011). The National Science Foundation now includes a cloud radar and lidar in its deployment pool, and has invested in them for its newest airborne platform for environmental research.
VOCALS-REx included the 94 GHz (3.2 mm wavelength) Wyoming Cloud Radar (WCR) and the 355-nm wavelength Wyoming Cloud Lidar (WCL; Wang et al., 2009) on the C-130 plane. The planners of VOCALS-REx also recognized the value of routine, ideally instantaneous, liquid water path (LWP) estimates. US aircraft missions have generally limited their use of microwave radiometry for boundary-layer cloud process studies, maybe because downward-looking radimeters need to account for surface reflectance and emission (e.g., Liu et al., 2003), or, because the instrument sizes can challenge payload restrictions. The search for a suitable radiometer ultimately led to a unique (for stratocumulus studies) feature of the VOCALS Regional Experiment: the inclusion of the millimeter-wave 183-GHz G-band Vapor Radiometer (GVR) (Pazmany, 2007) on the C-130 plane. The radiometer fit into a standard wing-mounted canister. Its small radiometer size is in part a virtue of the small wavelength (∼1 mm) compared to those from more standard microwave instruments (∼10 mm). An upward-pointing orientation facilitated integration with data from an upward-pointing radar and lidar, and allowed the retrievals to use the isotropically-emitting cold space as the boundary condition.

A surface-based version of the GVR was deployed at Barrow, Alaska (Racette et al., 2005; Pazmany, 2007; Cadeddu et al., 2007; Cimini et al., 2007, 2009), where the high sensitivity of 170–197 GHz radiation to water vapor and liquid water emission within a cold, dry climate (WVPs < 5 mm) could be used to full advantage. The airborne GVR has also been applied to wintertime mixed-phase clouds in the dry western United States (Wang et al., 2011). As shown later, the WVP values retrieved by the GVR for the free-troposphere of the southeast Pacific are similar to surface WVP values for Barrow, Alaska (see also the discussion of measurements from the Chilean altiplano in Turner and Mlawer, 2010). The low free-tropospheric WVP values reinforce the idea that a broad window in the infrared spectrum is available for stratocumulus cloud top radiative cooling in the southeast Pacific. The free-tropospheric WVP retrievals have also allowed an investigation of satellite microphysical retrieval sensitivity to above-cloud water vapor absorption for this region (Painemal and Zuidema, 2011).

For the lower-altitude sub-cloud legs, the water vapor overlying the plane fully saturated the center 183 GHz water vapor absorption line. The emission from the wings of the absorption line retained sensitivity to LWP, however, and LWP can be retrieved given an input water vapor path specified independently using other aircraft datasets. This is a new and different application of the GVR from its previous applications in the Arctic, and the GVR’s usefulness for cloud process studies within the dry tropics are undocumented. This is the purpose of this paper. The exercise will in the future contribute to the performance evaluation of another 183 GHz radiometer simultaneously deployed upon the Research Vessel Ronald Brown but of different design and calibration technique, the Radiometrics MP-183.

Specifically, we ask:

1. How well do GVR retrievals of free-troposphere WVPs compare to those from radiosondes?

2. Can stratocumulus LWPs be usefully retrieved from just the wing-line emission with other aircraft measurements specifying the boundary-layer WVPs, and if so, to what accuracy?

3. If so, what do the results tell us about the cloud adiabaticity?

The first question anticipates the use of the many radiosondes previously launched in this region. These radiosondes provide the inputs into forward calculations of brightness temperatures that are then used for the WVP retrievals, similar to Racette et al. (2005), Cadeddu et al. (2007), Cimini et al. (2007), and Cimini et al. (2009). The second question explores this study’s approach for retrieving LWP, of relying on aircraft state variables and the other cloud remote sensors to specify the boundary-layer WVP. Retrieval accuracy is primarily assessed by taking advantage of the pronounced cloud adiabaticity previously documented for this region (e.g., Bretherton et al., 2004; Zuidema et al., 2005), enough that adiabatic derivations have been used as a LWP proxy for VOCALS-REx (e.g., Bretherton et al., 2010). The GVR retrievals in turn are used to provide a larger statistical evaluation of cloud adiabaticity, establishing limits on the impact of entrainment and sub-cloud decoupling on cloud processes. This dual function was achieved by first selecting ideal adiabatic conditions for the retrieval assessment, and then evaluating adiabaticity statistics from a full dataset of approximately 60 000 1 Hz samples, equivalent to ∼6000 km of distance sampled at an aircraft flight speed of 100 m s⁻¹.

2 Background: initial retrieval design, further instrument description, and radiative absorption models

2.1 Initial retrieval considerations

Approaches for deriving LWPs from the GVR measurements include a neural net algorithm (Pazmany, 2007) and a four-channel physical retrieval (Cadeddu et al., 2007). These algorithms were developed from mid-latitude winter and Arctic temperature and moisture profiles that are not representative of subtropical stratocumulus conditions. In particular, the center absorption lines did not saturate. We developed initial expectations from forward radiative transfer calculations of brightness temperature based on radiosonde moisture and temperature profiles from seven buoy-tending cruises occurring between the years 2001 through 2008. These sampled the full diurnal cycle at regular intervals, and spent 6–7 days at 85° W, 20° S, and ∼3 days along 20° S to the coast.
(or vice versa). All but two of the cruises took place in October and all the post-2001 cruises used the more accurate Vaisala RS-90 series radiosondes. The cruises are described more fully in Zuidema et al. (2009), de Szoeke et al. (2010), de Szoeke et al. (2011) and references therein. Histograms of the water vapor paths for the years prior to VOCALS-REx and separately for the VOCALS cruise are shown in Fig. 1. Free-tropospheric WVPs (calculated from 1370 m altitude, the mean above-cloud aircraft height) during VOCALS-REx were low compared to other years, typically <0.5 mm with outliers up to three mm. Radiosonde-derived water vapor paths calculated at 150 m, a representative altitude for the aircraft sub-cloud legs, also show VOCALS-REx WVPs lower than for other years, with values between eight to twenty mm (Fig. 1b).

Several choices were made within the forward radiative transfer calculations based on the soundings. Brightness temperatures \( T_b \) calculations for the center lines can be sensitive to stratospheric water vapor (providing the basis for satellite remote sensing of moisture). The sensitivity is most pronounced for WVPs <1 mm and thus could potentially influence the free-tropospheric WVP retrievals. Racette et al. (2005) and Cadeddu et al. (2007) concluded that uncertainty in the specification of water vapor above 10 km could only explain a difference of ~1 K in the center line \( T_b \), however, later supported by Payne et al. (2008). Thus, a climatological stratospheric sounding extrapolated onto the radiosonde soundings reaching at least 50 hPa altitude was deemed sufficient, as also done within Cadeddu et al. (2007) and Payne et al. (2008). In addition, although in theory the \( T_b \) could be increased by scattering from high-altitude ice crystals, the observed cirrus fraction was low and clouds thin, with previous studies reporting negligible scattering for ice particle sizes <200 \( \mu \)m (e.g., Buehler et al., 2007; Cimini et al., 2007). We did not consider cirrus scattering. Vaisala RS-90 soundings do contain a small daytime solar heating bias, estimated at 5 % (Miloshevich et al., 2009), but otherwise brightness temperatures calculated from Vaisala RS-90 series radiosondes have compared well to GVR measurements in the Arctic (Payne et al., 2008) and to chilled mirror hygrometers (Miloshevich et al., 2006) using the above assumptions.

The downwelling \( T_b \)s corresponding to the 10th and 90th percentile of 300 Vaisala RS-92 radiosondes reaching 50 hPa are shown for the 0–220 GHz frequency range in Fig. 2, bottom panel, (black and red solid lines respectively), and for the 165–200 GHz range (top panel). The spectra corresponding to the above-cloud legs (purple and yellow lines) show the center lines retaining their sensitivity to moisture. At 150 m altitude the center absorption lines are fully saturated, while the wing lines retain a strong sensitivity to LWP regardless of WVP. The sensitivity of the two outer wing line \( T_b \)s to liquid water path is shown in more detail in Fig. 3. The \( \pm 7 \) GHz \( T_b \)s (red lines) are less sensitive to LWP changes in more moist conditions, while the outer wing lines (\( \pm 14 \) GHz, black lines) remain responsive to liquid water changes – \( \frac{\delta T_b}{\delta LWP} \) sensitivities between 0.6–1.5 K per 10 g m\(^{-2}\) depending on water vapor. Initial assessments suggested the \( \pm 7 \) and 14 GHz channels performed equally well. For the work presented here, the choice was made early on to base the LWP retrieval on just the \( \pm 14 \) GHz \( T_b \) because that radiance was more highly sensitive, and more linearly sensitive, to liquid water path than the \( \pm 7 \) GHz \( T_b \). Nevertheless, the combination of the \( \pm 7 \) and \( \pm 14 \) GHz \( T_b \)s could potentially reduce the sensitivity of the retrieval to uncertainties in the inputs, particularly the boundary layer WVP. Its assessment remains outside of the scope of the work presented here, but is a worthwhile future activity.

### 2.2 G-band vapor radiometer instrument design and calibration

The GVR measures brightness temperatures from four double-sideband channels, centered at \( \pm 1, \pm 3, \pm 7, \) and

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**Fig. 1.** (a) Histogram of free-tropospheric water vapor path values calculated from ship-board radiosondes prior to VOCALS-REx, and VOCALS-REx separately (black and red lines, respectively). Dotted lines indicate the cumulative frequency of occurrence. (b) Same as (a) but for the water vapor paths overlying the sub-cloud aircraft legs at 150 m altitude. All radiosondes were released between 0° and 23° S and east of 97° W and all reached at least 200 hPa.
Fig. 2. Downwelling brightness temperatures at 150 m altitude calculated using the 10th and 90th percentile of the 330 RS-92 radiosondes reaching 50 hPa (WVPs of 11 and 21 mm, black and red lines respectively), for 0–220 GHz (bottom panel) and 165–200 GHz (top panel). Cloudy-sky (LWP = 100 g m\(^{-2}\), cloud base at 1 km) spectra are shown as dashed lines. Above-cloud spectra at 1370 m are also shown (WVPs of 1.9 and 9.8 mm, purple and yellow lines, respectively). Thin vertical lines indicate GVR frequencies, these are combined into four double-sideband values within the radiometer.

±14 GHz off of the 183.31 GHz water vapor absorption line (Pazmany, 2007). The absorption spectrum for LWP is not symmetric around 183 GHz (evident in Fig. 3). The instrument bandwidths are 0.5, 1.0, 1.4 and 2.0 GHz for the ±(1, 3, 7, and 14) GHz center frequencies, whereas the \(T_b\) calculations were monochromatic and centered within the bandwidth. In addition, we assumed that each of the two sidebands contributed evenly to the measurement. All of these assumptions may affect the comparison of the GVR measurements to the radiosonde-calculated \(T_b\)s, but they seem reasonable given the passband frequency response shown in Fig. 2 of Pazmany (2007).

Salient details relevant for understanding the instrument calibration, provided within Pazmany (2007) and Cadeddu et al. (2007), are repeated here. The sky brightness temperatures are calculated internally using receiver gain and offsets determined from two reference loads at set temperatures (Pazmany, 2007); these were about 295 K and 333 K during VOCALS-REx. The receiver gain and offset were monitored through viewing the two loads approximately once every ten seconds on a continuous cycle. This frequent monitoring reduced the impact of receiver temperature fluctuations upon the calculated sky brightness temperatures. The absolute accuracy of the reported load temperatures is approximately 0.2 K (Pazmany, 2007), which could in theory correspond to an error in the sky brightness temperature of almost 3 K (Cadeddu et al., 2007). This is considered to be the most significant source of calibration error (Cadeddu et al., 2007).

An estimate of the absolute calibration was done in Barrow, Alaska by comparing the ground GVR measurements to radiosondes and other millimeter-wave radiometers and arrived at ∼2 K (Cadeddu et al., 2007; Cimini et al., 2009). An independent external assessment of the airborne GVR, done immediately prior to VOCALS-REx, found all four \(T_b\) measurements within 1 K of each other and of the calibration target. We assume a 2 K calibration error estimate here to be consistent with the stated accuracy of the internal temperature sensors (Pazmany, 2007).

The channel noise provides an estimate of the relative calibration (i.e., the accuracy of individual values compared to their neighbors). It depends on the integration time, bandwidth, and, receiver noise temperature, and is estimated at 0.2 K at a 200 ms integration time (Pazmany, 2007).
be assessed experimentally using a time series within highly homogeneous conditions. We evaluated several random 21-s time series from the outgoing ferry leg of flight RF9 at 5.3 km, and found representative standard deviations of (0.8, 0.3, 0.3, and 0.4) K for the ±(1, 3, 7, and 14) GHz lines. These support the published estimate and we chose 0.3 K as a representative value.

A temporal resolution of 1 s was desired for the reported brightness temperatures, to be on a common framework with that of other remote sensing measurements. This corresponds to an approximate 1 Hz measurement rate with occasional 0.25 s calibration gaps. The retrieved LWPs represent the average over a horizontal aircraft transect leg of ~30 m, estimated from the antenna beamwidth of 1.7° and an aircraft-distance-to-cloud of 500 m.

2.3 C-130 Flight details, Wyoming cloud radar, Wyoming cloud lidar, and ancillary datasets

Fourteen C-130 research flights sampled the SEP atmosphere in October and November of 2008, eight at night and six during the day. Four of the nighttime flights sampled along 20° S out to 85° W. Its flight plan, shown in Fig. 4, was typical of all missions in that it dedicated ten minutes each to sub-cloud legs (SC; ~150 m altitude), cloud-base legs, and above-cloud legs (AC). The above-cloud legs occurred at ~200 m above the cloud top for a mean above-cloud flight altitude of 1370 m, with cloud top heights varying between one to two km (Bretherton et al., 2010; Rahn and Garreaud, 2010). Other flights sampled open cellular convection with legs of 40 min duration. Several of the daytime flights sampled along the coast. More details on overall aircraft performance and individual flights are available within Wood et al. (2011b).

The WCR and WCL provided the cloud boundaries needed for the adiabatic LWP calculation. The radar was situated towards the rear of the plane and the lidar and GVR were located more centrally. For this study only the upward-pointing beam of the 94 GHz radar was used. The cloud radar detected cloud top given sufficiently strong scatterers, at a vertical resolution of 30 m. The cloud lidar possessed a vertical resolution of 3.75 m, and provided a sensitive detection of cloud base in the presence of precipitation (if more variable). At times, the presence of scud – a lower, secondary cloud layer formed from precipitation-moistened air – inflated the cloud thickness estimate attributed to the active overlying single-layer cloud (e.g., Fig. 14 of Bretherton et al., 2010). During the RF9 flight, the lidar computer clock occasionally drifted. This was corrected for as best possible by removing time discontinuities between the GVR and adiabatic LWP time series when they were expected to match. Excluding the RF9 flight, 60 590 1 Hz samples were collected.

The cloud base height (Z_{cb}) combined with the lifting condensation level was used as a measure of boundary layer mixing. (Z_{cb}-LCL) differences greater than 125 m were taken to indicate decoupled boundary layers, following Jones et al. (2011), who found the difference corresponded to an approximate 0.5 g kg^{-1} difference in the total mixing ratio between the bottom 25% and upper 25% of the boundary layer. Initial calculations found LCLs that exceeded Z_{cb}s, and an offset of 0.8 K was added to the reported dewpoint temperatures to bring the LCL and Z_{cb} into agreement for non-precipitating well-mixed conditions (Bretherton et al., 2010). This also allowed in-cloud relative humidity readings to reach 100 %. Last, the cloud base temperature (T_{cb}) was determined from a Heimann radiometric sensor measuring radiation between 9.5 to 11.5 μm with a stated accuracy of ±0.5 K.

Adiabatic calculations were only done when both radar and lidar identified their respective cloud boundaries, primarily selecting for thicker clouds. These comprised 40 % of the total samples, and the comparison of retrieved to adiabatic LWPs was restricted to these. For thin clouds the lidar could detect a cloud base but a radar-based algorithm might not identify a cloud top. For such cases, the altitude of the
radar return with the maximum reflectivity, if available, was used to establish the cloud top for the GVR LWP retrieval. Only time series containing a consistent lack of returns were deemed clear-sky. Statistics on the cloud boundaries for the sub-cloud legs only are summarized in Table 1.

2.4 Millimeter-wavelength absorption models

The correspondence between sonde-calculated \( T_b \)s and the GVR measurements also depended on the microwave absorption models used to convert radiosonde moisture and temperature profiles into microwave brightness temperature values. Previous work at 22 and 30 GHz frequencies favored the Liljegren et al. (2005) gaseous absorption model, itself modified from Rosenkranz (1998) for this stratocumulus region (Zuidema et al., 2005). For the current work we further modified the Liljegren et al. (2005) gaseous absorption model to include the water vapor continuum coefficients in Table IV of Turner et al. (2009). This was done to address concerns that LWPs retrieved using all four 183 GHz channels could be markedly smaller than those using the 22–30 GHz channels (Cadeddu et al., 2007), perhaps because of the spectroscopic specification of the 183 GHz absorption line (Payne et al., 2008) and/or the water vapor continuum line (Hewison, 2006). Our modification of the water vapor continuum coefficients brings intermodel differences to within 2 K over the window regions between 10 to 400 GHz Turner et al. (2009). The modification primarily impacts the\((14 \text{ and } 7) \text{ GHz channels.} \)

The GVR liquid water path calculations also relied on the Liebe et al. (1991) and Liebe et al. (1993) liquid dielectric models. Water permittivity models are more difficult to evaluate than gaseous absorption models, but the Liebe et al. (1993) model has produced results similar to that of other models in the 10°-20° C range (Cadeddu and Turner, 2011). The more important limitation to our radiative transfer calculations is an assumption of Rayleigh scattering and attenuation: at the radiometer wavelength of \( \sim 1.66 \text{ mm, Mie scattering and attenuation effects can be locally large within heavy stratocumulus precipitation. The scattering and attenuation effects do partially compensate for each other. We did not explore this further here, but this needs to be researched more thoroughly.} \)

3 Free troposphere water vapor path retrieval calculations and assessment

Measurements and calculations done for the above-cloud RF3AC1 leg are used to demonstrate the retrieval of the free troposphere WVP. Brightness temperatures calculated from 330 VOCALS-REx RS-92 radiosondes at 1335 m, the RF3AC1 aircraft altitude, are shown as a function of WVP in Fig. 5a. The \( T_b \) response to increasing WVP is mostly linear for the wing lines, while the center line shows a non-linear response as WVPs increase beyond 2 mm. The\( \pm 3 \text{ GHz} T_b \) is the most sensitive of the four lines to WVP variability between 1 to 3.5 mm, maintains the most linear response, and has the least \( T_b \) variability at a given WVP of the four channels. The spread in the \( T_b \) values at a given WVP indicates sensitivity to the water vapor vertical profile and its emission temperature at those heights (Racette et al., 2005). The soundings show that most of the meteorological \( T_b \) variability originated from free-tropospheric moisture variability at \( \sim 5 \text{ km and in the upper troposphere (Fig. 6).} \)

The off-shore free-tropospheric brightness temperatures did not typically vary much over the \( \sim 72 \text{ km distance of one leg (e.g., Fig. 5b).} \) This is expected for a strongly-subsiding atmosphere with little variation in the free-tropospheric WVPs. The average of nine RF3 above-cloud leg-mean \( T_b \) standard deviations (between-leg variation was little and AC1 was excluded) was \((3.2, 3.0, 1.8, 1.6) \text{ K for the } \pm(1, 3, 7, \text{ and } 14) \text{ GHz channels. The center line was the most variable as it was the most sensitive to low clear-sky WVP values. Near the coast, continental moisture outflow can give rise to gradients in the measured } T_b \text{s. The gradients were typically local, within } \sim 10 \text{ km (e.g., Fig. 5c).} \) We typically

<table>
<thead>
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<th>lidar returns (cloud base)</th>
<th>radar returns (cloud top)</th>
<th>adiabatic LWPs</th>
<th>clear-sky (no lidar/radar)</th>
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<tr>
<td>( N )</td>
<td>51 836</td>
<td>41 449</td>
<td>24 192</td>
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<tr>
<td>%</td>
<td>85.5</td>
<td>68.4</td>
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Fig. 5. (a) Simulated \( T_b \) for the \( \pm1, \pm3, \pm7, \text{ and } \pm14 \text{ GHz} \) GVR channels calculated for 330 RS92 soundings reaching 50 hPa (purple, blue, red, and black filled circles, respectively) at the RF3AC1 altitude of 1335 m, and GVR leg-mean values for RF3AC3 (blue diamonds, collected at 1548 m altitude and shown only for comparison) and 5-min mean values for RF3AC1 (yellow diamonds). Time series of GVR \( T_b \)s for (b) RF3AC3 and (c) RF3AC1, with \( \pm(1, 3, 7, \text{ and } 14) \text{ GHz} \) lines indicated in purple, blue, red, and black respectively.
estimated a leg-mean WVP value from a best-fit of the four GVR $T_b$s to the forward-calculated radiosonde-derived $T_b$s done at leg-mean aircraft altitude. Near the coast the leg measurements were split into two halves. Example WVP values are shown in Fig. 5a for the two legs shown in Fig. 5b and c. If an initial WVP estimate is made using a best-fit to the $\pm 3$ GHz line only, the deviations from the other channels suggest a conservative WVP error estimate of $\sim 0.2$ mm.

### 3.1 Comparison to R/V Ron Brown sonde

The above-cloud GVR $T_b$s were independently assessed by comparing GVR values from three C-130 overflights of the R/V Ronald Brown to $T_b$s calculated from ship soundings. The three overflights occurred on 25 October (RF5), 2 November (RF8), and 11 November (RF12). The comparisons are shown in Fig. 7. The nighttime 2 November comparison was the most robust (Fig. 7a), in that it was free of a known dry daytime humidity bias (Miloshevich et al., 2009), and a sounding was launched near in time to the overflight. The two aircraft above-cloud legs occurred 18 min prior and 41 min to a sounding, approximately 100 km east and 50 km west of the ship. The GVR $T_b$s exceeded the radiosonde-derived values, most notable for the center lines. The 11 November afternoon GVR $T_b$s also exceeded the radiosonde-calculated values, even after allowing for a 5% difference due to radiosonde solar absorption (Fig. 7b). The closest 25 October aircraft-ship match is also shown for completeness (Fig. 7c) but adds little information (GVR data were lacking from the direct overpass of the R/V Ron Brown).

The comparisons of the measurements to the radiosonde-calculated values (Fig. 7) suggest the GVR $T_b$s may be biased high. For 2 November, GVR-estimated WVPs of 0.0 mm and 1.5 mm (at 1.38 and 1.60 km altitude) exceeded the radiosonde-derived values of 1.42 and 1.27 mm by 0.6 and 0.2 mm respectively (Fig. 7a). The longitudinal $T_b$ gradient of the eastern flight leg, if extrapolated, may explain some of the 0.6 mm discrepancy. The comparison at $\pm 3$ GHz was the worst of the four channels, however, and can also be seen on Fig. 5. This line is nearest to the spectroscopic ‘pivot point’, at which specifications in line width and intensity should compensate for each other (Payne et al., 2008). Nevertheless, both Payne et al. (2008) and Turner et al. (2009) indicate that the spectroscopy of the center 183.31 GHz absorption line may require further refining, which is at least consistent with the offset shown here. We lacked the means for further evaluation. Almost all cases showed a good general correspondence, similar to that in Fig. 5. For above-cloud leg segments with little $T_b$ variability, we estimate the error in the retrieved WVP at 0.2 mm. The along-coast flights experienced the most $T_b$ variability and may have larger errors in their retrieved WVPs.
4 Cloud liquid water path retrieval

The free-tropospheric WVPs retrieved from the above-cloud legs were summed with boundary-layer WVPs computed from in-situ thermodynamic data, and the total WVP used as an input into the LWP retrieval. The boundary layer water vapor path was calculated, at 1Hz resolution, by integrating the in-situ water vapor mixing ratio from flight level up to the lifting condensation level. If the LCL was higher than half-way from the flight level to cloud base, the in-situ water vapor mixing ratio was integrated up to one-half the distance from the flight level to Zcb, and the cloud base mixing ratio used thereafter. This was done to account for decoupled conditions, for which the water vapor mixing ratio at cloud base might not match the in-situ vapor mixing ratio. The cloud base temperature (Tcb) was estimated from the larger of the Tcb estimated by the aircraft infrared radiometer, or, that estimated using a dry adiabatic lapse rate. Above cloud base, liquid water was adiabatically removed from the water vapor mixing ratio up to the radar-inferred cloud top.

Several limitations exist to this methodology. The radiometrically inferred Tcb can be underestimated if the cloud is in fact a greybody, i.e., when LWPs <30 g m⁻². The inferred Tcb can also be overestimated because of unaccounted-for water vapor emission. In highly-decoupled conditions, the use of the dry adiabatic lapse rate may generate an underestimate of the Tcb. An error in the cloud base temperature or height by 1 K or ~100 m is estimated to correspond to an error in the boundary-layer WVP of 0.13 mm, assuming a WVP of 20 mm distributed over 1.5 km depth. Similarly, the adiabatic assumption may overestimate the removed liquid water amount. An overestimate of, e.g., 50 g m⁻² liquid water will only underestimate the boundary layer water vapor path by 0.05 mm. Each of these errors is less than the 0.2 mm error ascribed to the free-tropospheric WVP. Although all of the limitations mentioned here are unfortunate, none are considered severe. This is ultimately confirmed by the good correspondence between the retrieved and adiabatic LWPs in situations where that is expected. In sum we assign a maximum error of 0.2 mm to the boundary layer WVP specification, with smaller values expected for well-mixed boundary layers.

The LWP retrieval itself was modified from that of Zuidema et al. (2005). An estimate of the cloud temperature (to which the millimeter-wave absorption is sensitive) was made from the cloud boundaries combined with temperature sounding composites done at three off-shore distances: coastal, 300–600 km, and 600–1000 km, similar to Fig. 6 but done for temperature. Errors of a couple of degrees in inferred cloud temperature only have a small impact on the millimeter-wave emission, and cloud boundaries were defined more leniently within the retrieved LWP calculation than for the adiabatic calculation. If local cloud boundary information was missing, either the neighboring or the leg-mean cloud boundaries were used. If no radar cloud top was determined during the entire leg (more likely to occur than no cloud bases), a cloud top height was estimated from the neighboring ascent and/or descent legs. For the RF9 flight (9 November), the cloud bases were solely estimated from the aircraft profile legs, and although GVR LWP retrievals were done, their accuracy is unknown. The retrieval then consisted of forward-computing brightness temperatures as a function of LWP, and iterating on the assumed LWP until a best-fit of the computed ±14 GHz Tb to the measured ±14 GHz Tb was found. In a couple of near-coastal cases, the LWP iteration could not always be made consistent with the composite temperature profile, adversely affecting the retrieval. Future work could be done to address this; RF13pocs2b may be one such case.

A summary of the major error sources within the GVR LWP retrieval is given in Table 2. Disregarded errors include Mie effects, daytime instrumental heating, monochromatic calculations, and the 50% sideband-weighting. Mie effects on both scattering and emission are considered the most significant of these errors. Errors will be higher within decoupled boundary layers, documented in Sect. 5, and near the coast, where the water vapor path overlying the plane is both higher and more variable. Errors in LWP are estimated from the sensitivities shown on Fig. 3. Figure 3 also shows that errors (sensitivities) decrease with increasing LWP and WVP.

5 Cloud liquid water path assessment

Three independent assessments of the GVR LWPs were done. One traditional approach evaluates the non-zero LWP retrievals in conditions deemed clear by the lidar (e.g., Turner et al., 2007). This is less fruitful for the persistently cloudy southeast Pacific conditions (Table 1), but was nevertheless applied to the 8% of the samples that qualified. We also evaluated the LWPs (or lack thereof) for thin clouds: cases where the lidar detected a cloud base, but the radar did not detect cloud. This occurred in almost 14% of the total samples and were further broken down by coupled and decoupled sub-cloud conditions (Zcb-LCL < or > 125 m respectively). The most statistically-robust approach compared the retrieved LWPs to the adiabatic LWPs (40% of the total samples). These were also divided further into well-coupled and decoupled cases.
at beginning and end of the time series. Pre-descent and post-ascent LWP values derived from the FSSP and radar-derived cloud top (blue).

5.1 Best-case comparison of GVR LWPs to adiabatic LWPs

Many sub-cloud legs took place in apparently well-mixed boundary layers, but most such layers possessed mean LWPs \( \sim 100 \text{ g m}^{-2} \) (for a mean cloud thickness \( \sim 300 \text{ m} \)). A cloud thickness uncertainty of 30 m, corresponding to the vertical resolution of the cloud radar, corresponds to a cloud liquid water path uncertainty of (30%,20%,15%) for a (200,300,400) m thick cloud. Thus, thick yet non-drizzling clouds within well-mixed boundary layers are ideal for assessing the GVR LWPs with their adiabatically-derived counterparts.

Our best-case example, from the 21 October RF3 flight, had a mean cloud thickness of \( \sim 450 \text{ m} \) and a mean LWP slightly above 200 g m\(^{-2}\) (RF3SC2, Fig. 8), with the cloud base and LCL often closely aligned. The above-cloud WVP was 1.2 and 1.5 mm before and after this leg; values were linearly interpolated across the sub-cloud leg. The column-maximum radar reflectivity fluctuated between -20 and -10 dBZ and a robust radar-inferred cloud top was observed continually. The bordering aircraft ascents and descents indicated well-mixed conditions with minimal cloud-top entrainment. The agreement between the GVR and adiabatically-derived LWP was excellent, with only a 2% difference between their means (Fig. 8). This was true even towards the end of the leg, where the boundary layer would be considered decoupled.

5.2 Statistical comparison of GVR LWPs to adiabatic LWPs

In intermittently-coupled legs the correspondence between the adiabatic and GVR LWPs was nevertheless often good (e.g., Fig. 8). The RF3 flight sampled a particularly deep boundary layer with high LWPs and significant precipitation, yet conditions remained overcast, inviting a “best-case” statistical comparison using data from the entire RF3 flight. The GVR and adiabatic LWP distributions from all nine RF3 sub-cloud legs are shown in Fig. 9. The two distributions were basically identical for LWPs of 200 and 400 g m\(^{-2}\), lending further confidence to the GVR retrievals. The GVR LWPs exceeded the adiabatic LWPs for LWPs \( > 400 \text{ g m}^{-2} \). One plausible explanation is that Mie scattering from the larger drop sizes encouraged an overestimate of the actual LWP. For LWPs \( < 200 \text{ g m}^{-2} \), the most likely explanation for the greater frequency of GVR-retrieved LWPs is the generally lower radar reflectivities of the thinner clouds, so that a cloud top height was less likely to be reported and an adiabatic calculation not done.

5.3 LWPs in clear-sky conditions

Skies considered clear by the lidar were likely to also be considered clear by the GVR. LWPs retrieved within the clear-sky conditions of the sub-cloud legs for thirteen research flights totaled 4918 samples, or 8.1% of the total (Table 1). Their distribution is shown in Fig. 10a. Clear skies, within this data sample, possessed a LWP cloud fraction of 3%, and 80% of the retrieved LWPs were \( < 20 \text{ g m}^{-2} \).
5.4 Cloudy skies without liquid

Skies for which the lidar sensed a cloud base but the radar did not determine a cloud top should correspond to thin, low-LWP clouds. These totaled 22% of the 60590 cases. The GVR-retrieved LWP were also broken down into coupled and decoupled conditions (N=6179 and 7100 respectively, or 11% and 12.5%). The distributions of the LWP, when they were retrieved, are shown in Figs. 10b and c. The LWP retrieved within decoupled conditions were generally lower than in coupled conditions. This is physically plausible because upward moisture transport will be discouraged by the decoupling, but can also be a retrieval artifact. In decoupled conditions the water mixing ratio near cloud base may be reduced from that used to estimate the boundary-layer WVP, and the retrieved LWP will be underestimated or not retrieved at all. This appears to occur in some cases, as in 74% of the decoupled conditions, the GVR retrieval was likely to miss the cloud, compared to 59% of the coupled conditions. Nevertheless, reassuringly, the LWP, when retrieved, were low, mostly < 40 g m^{-2}, in both coupled and decoupled conditions, as expected.

6 Project-summary results

6.1 Free-tropospheric WVPs

The distribution of the retrieved leg-mean free-tropospheric (above-cloud) water vapor paths are shown as a function of longitude for all fourteen research flights in Fig. 11a, and as a function of latitude for the near-coastal legs (Fig. 11b). Free-tropospheric WVPs were often below 2 mm, and occasionally higher near the coast, particularly above the Arica Bight. Two of the WVP values approaching 10 mm came from the first and last above-cloud leg for the 23 October (RF4) flight, near Arica. Flight notes reported a haze and pollution layer at 3 km, with the aircraft sounding out of Arica showing relative humidities > 30% above 3 km. The other high free-tropospheric WVP value of 8 mm came from 2 November (RF8) from the first above-cloud leg, also near Arica, after the aircraft ascent had sampled a second moisture layer between 2–3 km, with RH ~ 30%, above the boundary layer. Most of the higher-WVP points west of 75° W came from the daytime flights.

Sounding-derived water vapor paths are also shown in Fig. 11 for VOCALS-REx (orange, 207 sondes) and the October–November cruises in 2001, 2003, 2006 and 2007 (purple, 293 sondes), calculated using the mean above-cloud aircraft altitude for each 2.5° bin. These show that the GVR- and sonde-derived values for 2008 were similar, despite different sampling days, with a drier free-troposphere in 2008 than in other years (consistent with Fig. 1 but using the GVR-derived dataset).
particular feature of the southeast Pacific divided into three longitude bands for the four nighttime flights dedicated to a 20° S nighttime flight pattern out to 85° W (21, 23, 25 October and 6 November; RF 3, 4, 5 and 10) subdivided by longitude; (a) 80° W–86° W, (b) 75° W–80° W, and (c) 70° W–75° W. LWP bins of 20 g m$^{-2}$. The cumulative frequencies for both the GVR-retrieved (black) and adiabatic (red) LWP correspond to the right y-axis.

6.2 LWPs

As shown in Fig. 9, the GVR LWP sample size for RF3 was larger ($N = 5471$) than the adiabatic sample size of 4003, because of the GVR’s ability to retrieve LWPs for thin clouds lacking radar-determined cloud tops. Most of the thin clouds occurred near the coast. This is shown by the summary statistics divided into three longitude bands for the four nighttime flights that flew along 20° S out to 85° W and back on 21, 23 and 25 October, and 6 November (RF 3, 4, 5, and 10) in Fig. 12. These flights flew the pattern shown in Fig. 4, leaving Arica, Chile near midnight and returning at dawn. Near the coast 45% and 86% of the GVR LWPs were <40 and 100 g m$^{-2}$ respectively. Mean±standard deviation (median) LWPs between 70°–75° W were 67±44 (61) g m$^{-2}$. The near-midnight sampled value appears representative of the diurnal-mean satellite-derived values shown in O’Dell et al. (2008) and O’Neill et al. (2011).

Further west, the mean (median) LWPs increased to 123±87 (105) g m$^{-2}$ between 75°–80° W, and to 166±152 (120) g m$^{-2}$ between 80°–86° W. Higher LWPs west of 80° W are consistent with deeper boundary layers and a greater propensity for precipitation (Leon et al., 2008; Zuidema et al., 2009; Bretherton et al., 2010), as well as a sampling time approaching a near-dawn diurnal maximum (Zuidema et al., 2005; O’Dell et al., 2008; O’Neill et al., 2011), but also here 23% and 50% of the GVR LWPs <40 and 100 g m$^{-2}$. As already shown in Fig. 9, clouds with retrieved LWPs between 100 and 400 g m$^{-2}$ matched adiabatically-calculated LWP values exceedingly well. This is consistent with previous results based on a smaller sample size (Zuidema et al., 2005), but not with North Atlantic observations (e.g., Pawlowska and Brenguier, 2000), and a high adiabatic fraction (LWP$_{\text{retrieved}}$/LWP$_{\text{adiabatic}}$) may be a particular feature of the southeast Pacific stratocumulus.

The distribution of retrieved and adiabatic LWPs from the subcloud legs for all research flights except RF9 is shown in Fig. 13, subdivided by coupling. Conditions were approximately evenly divided between coupled and decoupled: 52% and 48% (Table 3). The retrieved LWP cloud fraction was 71% for the well-coupled boundary layers, versus 67% for the decoupled boundary layers, with slightly more low-LWP values retrieved within coupled conditions (65% of the LWPs were <100 g m$^{-2}$ versus 61% for the decoupled samples). The adiabatic LWP cloud fraction, in contrast, was higher for the decoupled cases (51%) than for the well-mixed cases (42%). This can be explained by more radar cloud top returns for the decoupled cases because of precipitation within open cellular convection (Fig. 13c). The distributions of the GVR and adiabatic LWPs under decoupled conditions were also more similar than under coupled conditions, possibly because precipitation decoupling a boundary layer thermodynamically, may also reflect a stronger circulation that is otherwise encouraging mixing (see, e.g., Fig. 8). Overall, liquid water paths greater than ∼160 g m$^{-2}$ were determined well by both adiabatic and millimeter-wave approaches under both coupled and decoupled conditions, and the biggest difference between the coupled and decoupled cases was in the detection of thin clouds.

The liquid water path frequencies as a function of the independently-derived cloud thickness are shown in Fig. 14. Cloud thicknesses, available for all adiabatically-calculated LWPs (40% of the total samples; see Table 1), were extended a further 23% by establishing the altitude of the maximum dBZ return as cloud top for those samples with missing cloud tops. The adiabatic LWPs are more tightly constrained by the cloud thickness values than the retrieved LWP values, as would be expected. In this figure, one-half of all the GVR-retrieved LWPs correspond to clouds with recorded thicknesses <180 m thick, and this excludes 15% of the clouds with retrieved LWPs and no indicated cloud top.
Table 3. Number of 1 Hz samples contributing to the distributions shown in Fig. 13. These represent all samples from thirteen research flights with well-defined cloud bases and LCLs.

<table>
<thead>
<tr>
<th></th>
<th>N</th>
<th>N_{hwp}</th>
<th>N_{adiab}</th>
</tr>
</thead>
<tbody>
<tr>
<td>coupled</td>
<td>24 994</td>
<td>17 817</td>
<td>10 401</td>
</tr>
<tr>
<td>decoupled</td>
<td>26 842</td>
<td>18 048</td>
<td>13 791</td>
</tr>
<tr>
<td>total</td>
<td>51 836</td>
<td>35 865</td>
<td>24 192</td>
</tr>
</tbody>
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Figure 14 allows us to conclude that the thinner (LWP < 100 g m\(^{-2}\)) clouds are also adiabatic. Histograms of the adiabatic fraction, or ratio of retrieved LWP to the adiabatically-calculated value, are shown more explicitly for coupled and decoupled boundary layers in Fig. 15. For well-mixed boundary layers, the mean (median) adiabatic fraction is 0.81 (0.86) while for decoupled boundary layers the adiabatic fraction is 0.85 (0.96) (after disallowing adiabatic fractions > 1.0).

7 Conclusions

Brightness temperatures from a millimeter-wave radiometer on board the NCAR C-130 plane were used to retrieve above-cloud water vapor paths, and, were combined with remotely-sensed cloud boundaries and in-situ thermodynamic data to retrieve liquid water paths during the sub-cloud and cloud base legs. A particular strength of the higher 183±14 GHz frequency is the ability to retrieve low liquid water paths. The leg-mean LWP error in well-coupled conditions can originate from incorrect specification of the above-cloud WVP or boundary layer WVP (each estimated at 0.2 mm, or 6–10 g m\(^{-2}\)), and a 2 K absolute and 0.3 K relative uncertainty in the GVR \(T_b\), corresponding to 10–20 g m\(^{-2}\) and 2-3 g m\(^{-2}\), respectively (Table 2). Based on these numbers we conservatively estimate the absolute leg-mean LWP uncertainty to be ~20 g m\(^{-2}\) and the relative (within-leg) uncertainty at ~3 g m\(^{-2}\) in well-mixed conditions. In decoupled conditions, for which the boundary layer WVP specification is less certain, we estimate a higher absolute leg-mean LWP uncertainty of 25–30 g m\(^{-2}\). The uncertainties will be higher in heavily-precipitating conditions. Further work could be done to improve LWP retrievals in decoupled conditions within open cellular convection using all four GVR channels from the cloud-base legs.

The retrieved above-cloud water vapor paths were on par with surface-based Arctic values, further confirming the presence of a very dry troposphere in the southeast Pacific.
Both boundary-layer and free-tropospheric air east of 75° W are likely to encounter continental aerosol (Bretherton et al., 2010). The thin, adiabatic clouds are likely to be particularly susceptible to aerosol (e.g., Petters et al., 2011; Kim et al., 2008, and references therein), which may help explain the global-maximum albedo susceptibilities derived by satellite for the southeast Pacific (Oreopoulos and Platnick, 2008). Quantitative observational estimates of the first aerosol indirect effect can vary greatly, for reasons at least partially attributable to scale and to the choice of liquid water path as a constraint (McComiskey and Feingold, 2011). The complementary airborne cloud measurements at process-level scales may help estimates of the first aerosol indirect effect made using different platforms converge for this region, and have been used to understand precipitation susceptibility for this region (Terai et al., 2011).

Data from the three cloud remote sensors have been combined into Integrated Datasets publicly available for each research flight through the NCAR Earth Observing Laboratory data archive (http://data.eol.ucar.edu/master_list/?project=VOCALS), with plots for each research flight available through http://www.rsmas.miami.edu/users/pzuidema. The combined datasets free users from obtaining data for the individual instruments and from some of the data processing. The VOCALS experiment was the first and last field deployment on the C-130 of the airborne GVR, but the GVR has features that recommend further aircraft deployments of similar instrumentation. It is small, easily integrated onto an aircraft platform, inexpensive, and its upward-pointing orientation is conducive for integration with other datasets including concurrent sub-cloud air sampling, as well as providing a cold space background for the LWP retrievals. The retrievals could in theory be done instantaneously through a neural net (Pazmany, 2007). The millimeter-wavelengths are not optimal for the subtropical stratocumulus region, but could be combined with lower-frequencies focused on establishing the free-troposphere sampled during VOCALS was drier than in other years, with values ranging from one to two mm off-shore to almost one cm near the coast.

Clouds with LWPs <100 (40) g m\(^{-2}\) composed almost two-thirds (one-third) of the entire dataset. Thin clouds were more frequent near the coast, with a mean (median) LWP of only 67 (61) g m\(^{-2}\) between 70°–75° W. These values were measured near midnight and seem representative of their diurnal-mean (O’Dell et al., 2008; O’Neill et al., 2011). The GVR-retrieved LWPs compared well against adiabatically-derived LWPs at almost all LWP values (see Fig. 14 in particular), with 50% of the adiabatic fractions (LWP\(_{\text{retrieved}}\)/LWP\(_{\text{adiabatic}}\), cloud thicknesses >100 m only) exceeding 0.86. The high adiabatic fractions further support a view of southeast Pacific stratocumulus as occurring under strong, sharp inversions (e.g., Bretherton et al., 2010), with circulations strong enough to sustain cloud even in the presence of surface precipitation approaching 10–20 mm day\(^{-2}\) (Wood et al., 2011a), and with less cloud-top dilution than is observed, for example, in North Atlantic stratocumulus (e.g., Pawlowska and Brenguier, 2000).

Fig. 15. Distribution of adiabatic fraction (LWP\(_{\text{retrieved}}\)/LWP\(_{\text{adiabatic}}\)) for (a) coupled and (b) decoupled boundary layers. Only clouds with thicknesses exceeding 100 m were included, for a total N of 22 718.

References

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